

# *charmonium 'spectroscopy' from lattice QCD*

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will cover:

- \* sub-threshold charmonium
- \* excited states in charmonium
- \* radiative charmonium physics

displays most of the 'fundamental'  
problems  
pions show up less

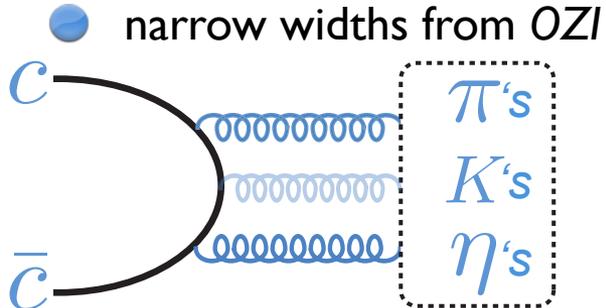
won't cover:

- \* bottomonium  
please address questions to  
NRQCD experts (I'm not one)
- \* heavy-light  
light-quark dominated?  
pions
- \* baryons

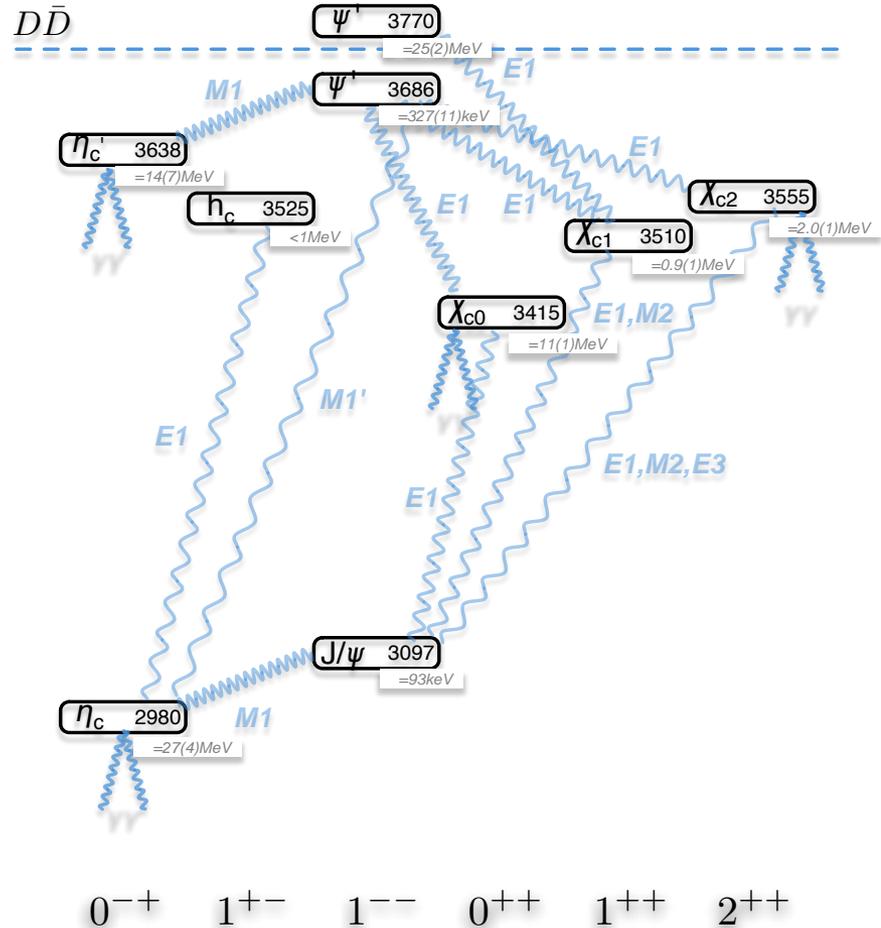
personal opinions - will find out in the next 30mins if these are controversial

# spectrum situation

- will focus on charmonium since that's where a lot of the action is currently
- sub-threshold states now mostly established



- radiative transitions
  - e.g.  $\chi_{c0} \rightarrow J/\psi \gamma$
- $\pi\pi$  transitions
  - e.g.  $\psi(3686) \rightarrow J/\psi \pi\pi$
- two photon decays
  - e.g.  $\chi_{c2} \rightarrow \gamma\gamma$



# two-point functions

- basic spectral information comes from *hadronic two-point functions*

$$\begin{aligned} C_{ij}(t) &= \sum_{\vec{x}} \langle 0 | \mathcal{O}_i(\vec{x}, t) \mathcal{O}(\vec{0}, 0) | 0 \rangle \\ &= \sum_N \langle 0 | e^{Ht} \mathcal{O}_i(\vec{0}, 0) e^{-Ht} | N \rangle \langle N | \mathcal{O}_j(\vec{0}, 0) | 0 \rangle \\ &= \sum_N \langle 0 | \mathcal{O}_i(\vec{0}, 0) | N \rangle \langle N | \mathcal{O}_j(\vec{0}, 0) | 0 \rangle e^{-E_N t} \end{aligned}$$

- the states labeled by **N** are **all** the eigenstates of the QCD Hamiltonian with the external quantum numbers of  $\mathcal{O}$
- the interpolating fields (or operators)  $\mathcal{O}$  are combinations of quark and gluon fields
  - e.g. a simple (local) interpolating field for a pseudoscalar is  $\bar{\psi}(\vec{0}, 0) \gamma^5 \psi(\vec{0}, 0)$
  - if the quark fields correspond to charm quarks, the sum over **N** includes all pseudoscalar states that have a non-zero amplitude to be in a  $\bar{c}c$  Fock state with both at the origin

# sub-threshold

- sub-threshold states are ideal for lattice computation (‘gold-plated’)
- neglect the *OZI* suppressed decays - i.e. don’t compute disconnected correlators

● e.g. using fermion bilinears  $C(t) = \langle \bar{q}(t)\Gamma q(t) \cdot \bar{q}(0)\Gamma q(0) \rangle$

$$= \sum_{\{U\}} Q_{t,0}^{-1}[U]\Gamma Q_{0,t}^{-1}[U]\Gamma \quad + \quad \sum_{\{U\}} Q_{t,t}^{-1}[U]\Gamma Q_{0,0}^{-1}[U]\Gamma$$


The diagram illustrates the decomposition of the correlator into connected and disconnected parts. On the left, a single loop with vertices labeled  $t$  and  $0$  is shown, labeled "connected". On the right, two separate loops are shown: one with vertices labeled  $t$  and  $t$ , and another with vertices labeled  $0$  and  $0$ , labeled "disconnected".

- most obvious problem in the past is the hyperfine structure, especially the  $J/\psi$ - $\eta_c$  splitting
  - believed to be a short distance effect, so must have short distance physics ‘right’
  - lattice is discrete on a scale  $a$
  - physics at shorter distance scales controlled by the particular discretisation used - “improved actions”

# sub-threshold

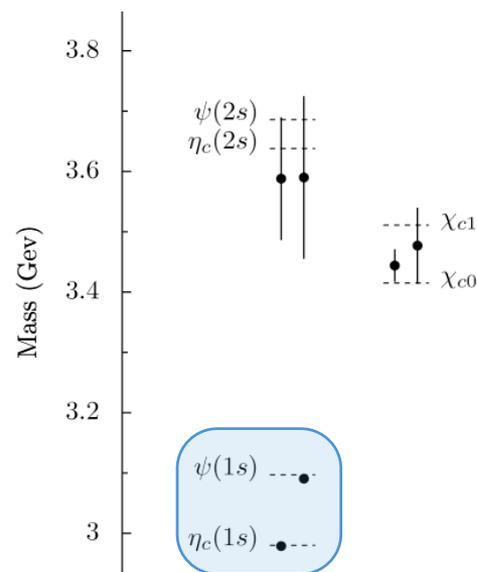
- also important that the coupling constant is right at short distances
  - usually set at large distances  $\rightarrow$  runs  $\rightarrow$  need right  $\beta$  function
  - so quenching the light quarks is a bad idea

$$g^2(k) = \frac{g^2}{1 + \frac{g^2}{3(4\pi)^2} (33 - 2N_f) \log \frac{k^2}{\Lambda^2}}$$

- one recent example using an improved action and dynamical light quarks

- *Phys.Rev.D75:054502,2007 (HPQCD & UKQCD)*
- $a \sim 0.09 \text{ fm}$  ,  $m_\pi \sim 250 \text{ MeV}$
- no disconnected correlators

- disconnected diagrams remain the largest challenge (both computationally and theoretically) in this region (*ask me - I have more*)



# ***fine structure***

- newly observed spin singlets,  $\eta_c(2S)$ ,  $h_c$ , are of interest
- especially their splittings from the  $\psi(2S)$ ,  $\chi_{cJ}$ 
  - *(a good idea to use a basis of operators - technical analysis issue)*
  - this sort of precision measurement requires improved actions, probably light sea quarks and probably inclusion of disconnected diagrams
  - I defer to precision experts here...
    - improving the precision is a very important problem
    - I'd like to focus on the 'other' problem - what further quantities can be computed (worry about the precision later - soon hopefully)

# excited states

- lightest state with each  $J^{PC}$  is technically easy to extract  $C(t \rightarrow \infty) \rightarrow e^{-m_0 t}$
- excited states are more challenging
  - technical problem / theoretical problem

## **technical problem:**

how do we actually extract reliable estimates of spectral quantities from correlators?

I believe this is 'solved':

**variational solution to  
matrix of correlators**

## **theoretical problem**

what happens to the spectrum of states when they can decay?  
i.e. "how does one deal with resonances?"

'Realistic' investigation of this has just begun:

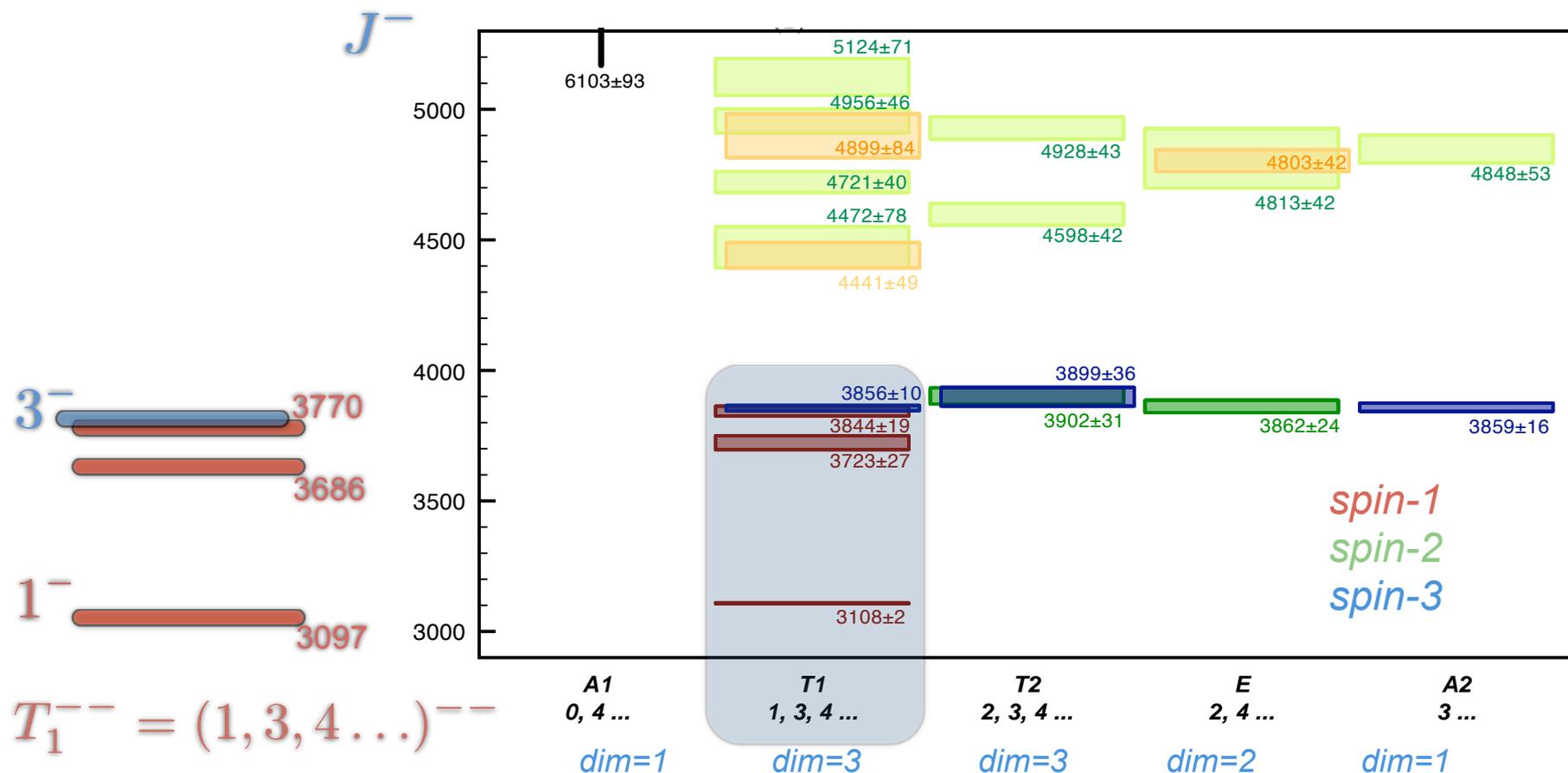
**Lüscher Method  
spectrum in a finite box  $\rightarrow$  S-matrix**

# excited states - technical problem

- compute a matrix of correlators  $C_{ij}(t) = \langle \mathcal{O}_i(t) \mathcal{O}_j(0) \rangle$
- using a bunch of interpolating fields, e.g.  $\bar{\psi} \Gamma \psi$     $\bar{\psi} \Gamma \overleftrightarrow{D}_k \psi$     $\bar{\psi} \Gamma \overleftrightarrow{D}_j \overleftrightarrow{D}_k \psi$
- a variational solution can be found - utilises the orthogonality of state vectors
- very successful at extracting multiple excited states in a given  $J^{PC}$

# variational method

e.g. recent (**quenched**) “charmonium” study (vector channel)



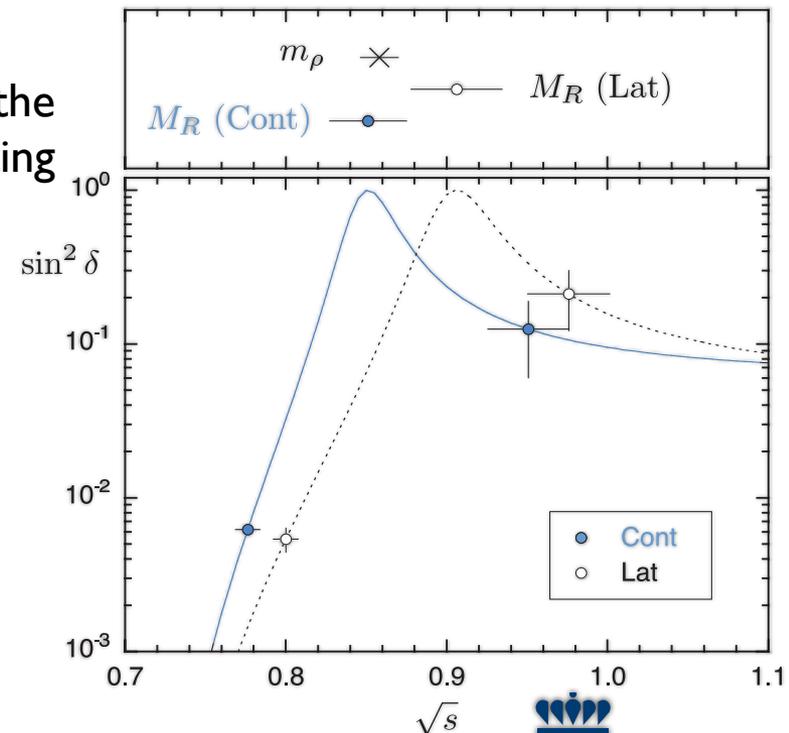
to my knowledge, the first time this has been seen

# excited states - theoretical problem

- how do we extract the mass and width of a resonance?
  - e.g. consider the  $\rho$  meson which decays to  $\pi\pi$
  - might try  $C(t) = \langle \bar{q}(t)\gamma^i q(t) \cdot \bar{q}(0)\gamma^i q(0) \rangle \rightarrow e^{-E_0 t}$  as  $t \rightarrow \infty$
  - will give us the lightest  $I^{G}$  eigenstate of QCD
    - this is two pions in a  $P$ -wave - rest energy =  $2m_\pi$
    - in an infinite box there are a continuum of such states
    - looks hopeless
- Lüscher (and others) have shown that in the **finite box** we work with in lattice QCD, there is a mapping between the energy levels extracted and the elastic scattering matrix (i.e. the phase shift)
  - in a periodic finite box, all energy levels are discrete so perhaps we can extract a small number and infer something about resonances?
  - one recent 'successful' application

# $\rho$ meson as a $\pi\pi$ resonance

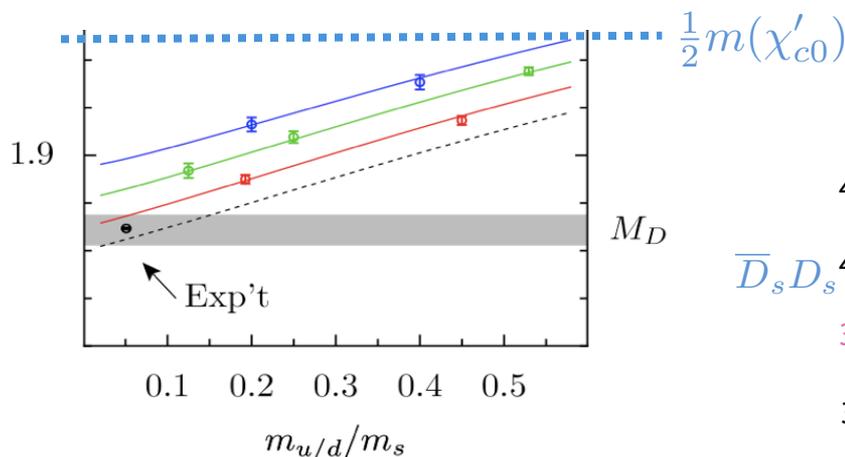
- CP-PACS recently used the Gottlieb-Rummukainen extension of the Lüscher method to study the  $\rho$  meson as a resonance in  $\pi\pi$
- at a fixed lattice volume they extracted two energy levels using a correlation matrix constructed from
  - a “ $\rho$ -like”  $\bar{q}q$  (wavefunction at the origin) operator
  - a “ $\pi\pi$ -like” separated  $\bar{q}q - \bar{q}q$  operator
- by comparing the extracted energy levels to the expected discrete levels for two non-interacting pions they inferred the  $\pi\pi$  phase-shift
  - simulation not at physical point
    - $m_\rho/m_\pi \sim 2.4$
    - $a \sim 0.2$  fm
    - $Lm_\pi \sim 4$



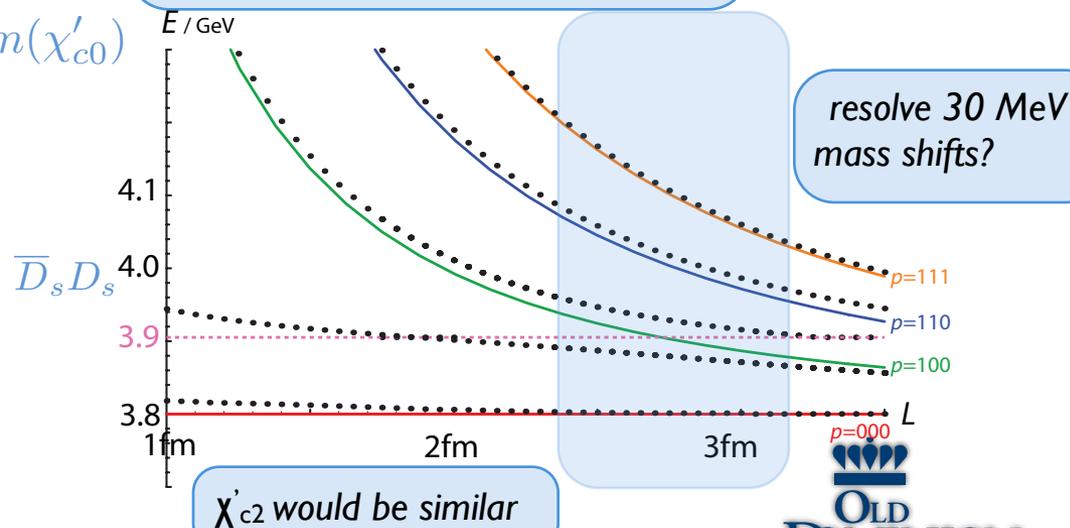
# charmonium application?

- the first excited  $\chi'_{c0}$  state would be a good place to try this out in charmonium
- observed  $\chi'_{c2}(3930)$  suggests  $\chi'_{c0}(\sim 3900)$  for which only  $\bar{D}D$  is open
  - is this practical on current lattices?
  - depends upon the  $D$ -meson mass at available quark masses
  - and the available volumes
  - and the energy resolution

HISQ on MILC



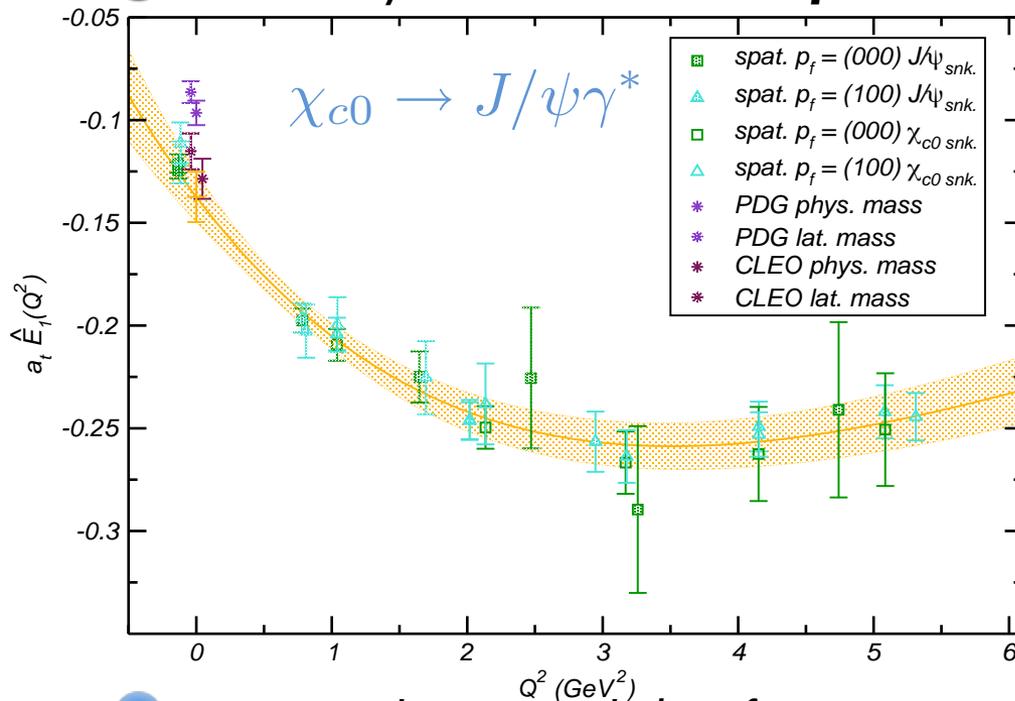
assuming  $m_D \sim 1.9\text{GeV}$ ,  
and  $m(\chi'_{c0}) \sim 3.9\text{GeV}$ ,  $\Gamma(\chi'_{c0}) \sim 50\text{MeV}$   
would get something like



# other 'spectroscopic' quantities

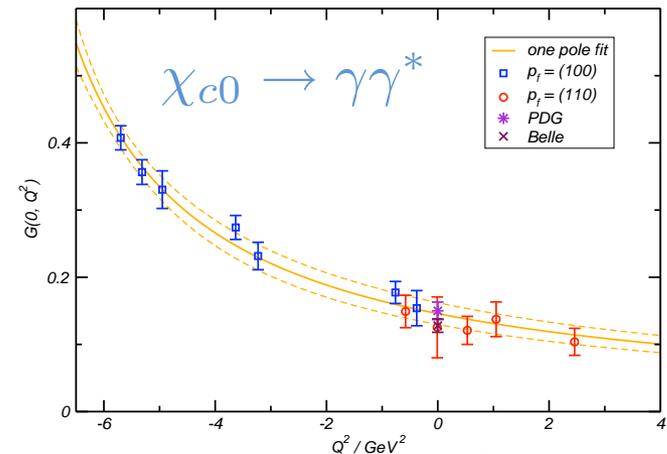
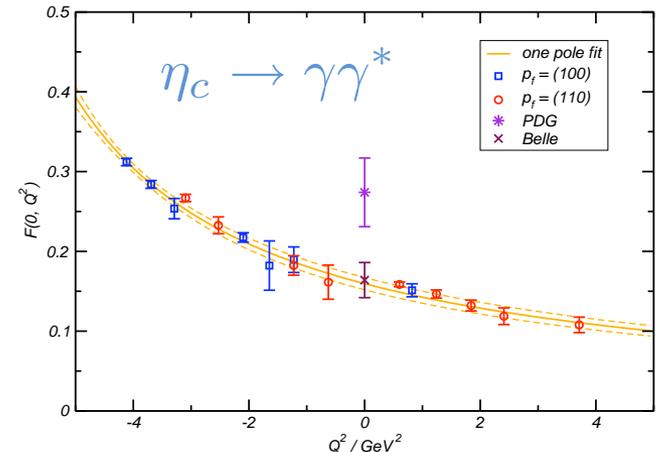
- radiative transitions and two-photon couplings can be obtained from three-point correlators
- N.B. two-photon fusion is the production mechanism for the new  $\chi'_{c2}(3930)$

so far only calculations are on **quenched** lattices



no excited states or  $J > 1$  so far

JLab group working on it - *much* harder



# ***X(3872) - an interesting case***

- *X(3872)*
  - unreasonably close to thresholds for  $D^+D^{*-}$ ,  $D^0\bar{D}^{*0}$  (1 MeV away?)
  - so close that isospin violation comes into play
- **no hope** of tuning everything to the precision required for direct lattice study
  - if 'binding' energy really  $\pm 1$  MeV, potentially long distance tail to wavefunction
- but still interesting to observe behaviour as quark mass & lattice volume change
- perhaps as input to effective field theory models?

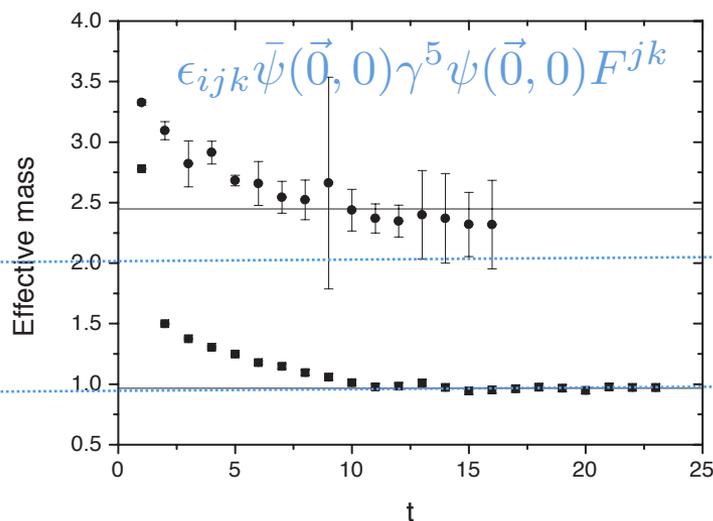
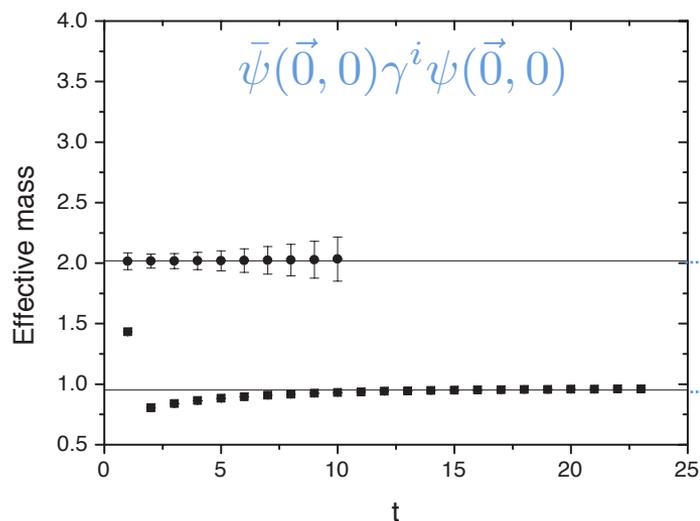
# $Y(4260)$ ( $1^{--}$ )

- model suggestions that it might be a non-exotic hybrid
- on a lattice this guy will be very tough!
  - above threshold for decay to  $\bar{D}D$ ,  $\bar{D}D^*$ ,  $\bar{D}^*D^*$ ,  $\bar{D}_sD_s$ ,  $\bar{D}_sD_s^*$
  - and above the  $J/\psi$ ,  $\psi(2S)$  and probably three resonances,  $\psi(3770)$ ,  $\psi(4040)$ ,  $\psi(4150)$
- need some major theoretical advances to consider this state
  - but far from understood phenomenologically!
    - *decay into  $\pi\pi J/\psi$  seen, but hadronic width is such it must be decaying elsewhere too*
    - *but doesn't show up as a peak in the exclusive*
    - *or (visibly) in the new exclusive data (CLEO, Belle, BaBar)*
    - *needs a (unitary!) coupled-channels fit including interferences properly*

modelling job

# interpolating fields & state interpretation

- it is a common mistake (even one made by some lattice theorists) to think that the “internal” structure of the interpolating field directly tells you the nature of states
  - e.g. any state produced using  $\bar{\psi}(\vec{0}, 0)\gamma^i\psi(\vec{0}, 0)$  must be a “conventional”  $\bar{c}c$   $1^{--}$
  - any state produced using, say,  $\epsilon_{ijk}\bar{\psi}(\vec{0}, 0)\gamma^5\psi(\vec{0}, 0)F^{jk}$ , must be a “hybrid” because of the gluonic factor
  - see e.g. “*Gluonic excitation of non-exotic hybrid charmonium from lattice QCD*”  
X-Q. Luo & Y. Liu (PRD74 034502)*quenched*



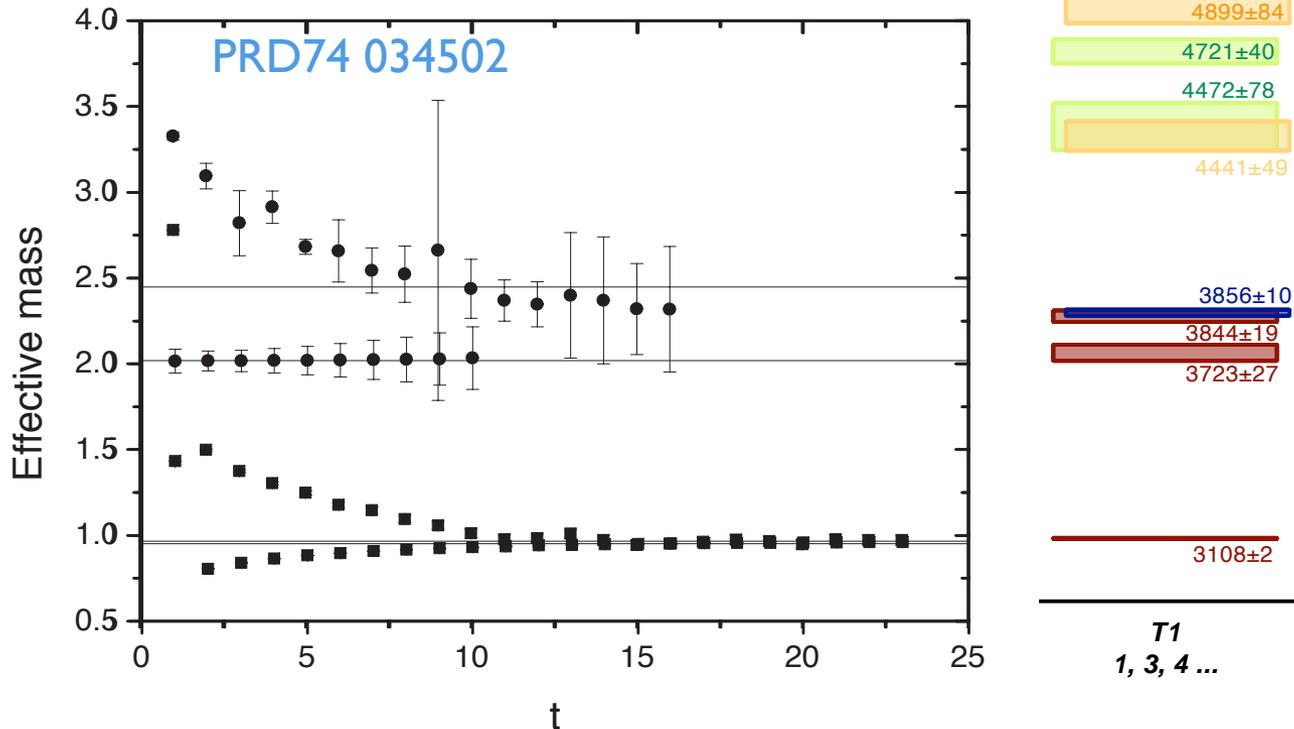
claim a hybrid !

$J/\psi$

if a “hybrid” is anything overlapping the second operator, then do we have to rethink the  $J/\psi$  ?

# large basis of operators

- similar (quenched study) using a matrix of correlators (including comparable operators)

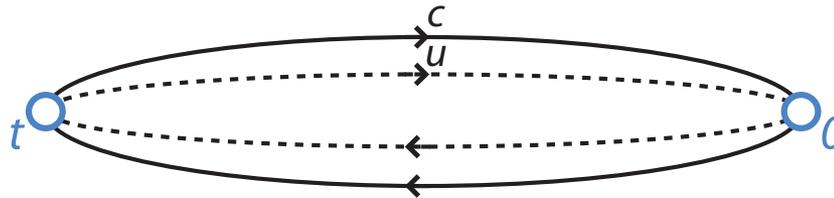


- appears to be far more mundane explanations for these correlators

# more (quenched) $Y(4260)$ claims

- Chiu & Hsieh “ $Y(4260)$  on the lattice” (PRD73 094510)
- compute correlators with multi-quark operators, e.g.  $(\bar{q}\gamma_i c)(\bar{c}q)$
- all quarks at same space point

- connected diagrams only

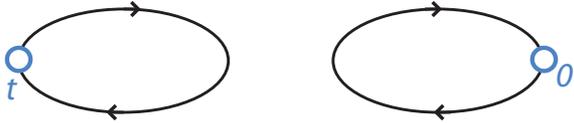


- technically this means  $isospin=1$ , but  $isospin=0$  might be degenerate?
- find ground state masses in the region of 4300-4500 MeV (after a crude extrapolation in  $m_q$ )

# $Z^+(4430) ??$

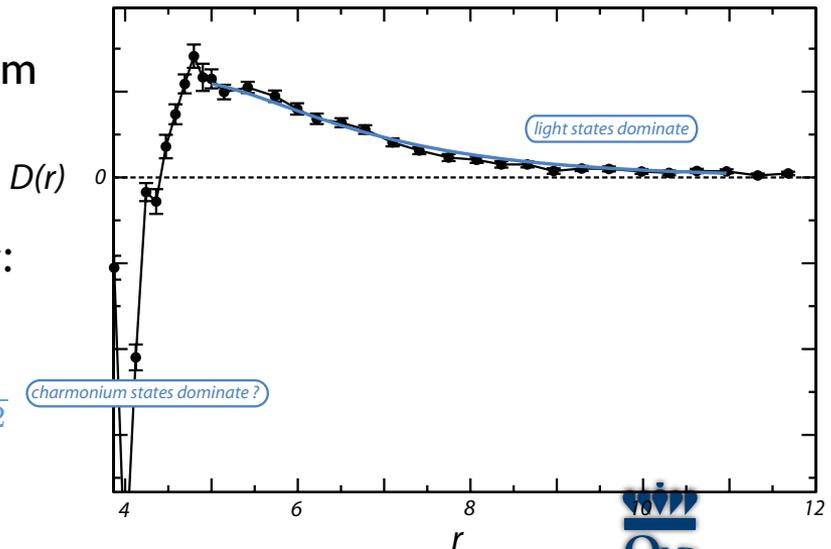
- 'seen' as a resonance in  $\pi^+\psi(2S)$  by Belle via  $B$ -decays
- isospin one with large affinity for charmonium - might be interesting !?
- potentially lots of decay channels open so similar problem to  $Y(4260)$

# disconnected contributions



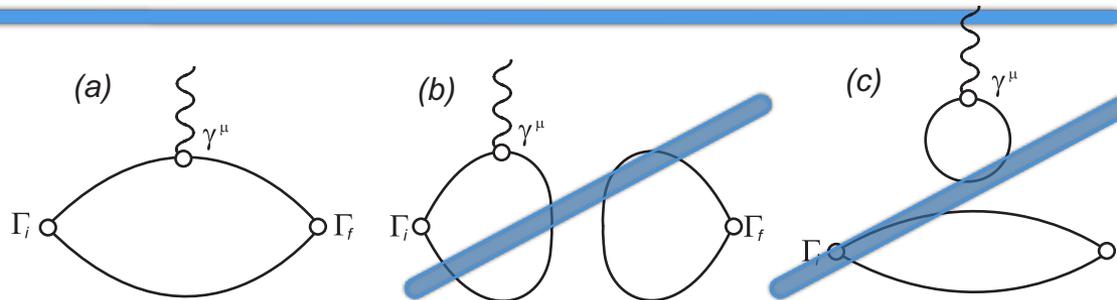
- technically challenging (noisy signals)
- theoretically challenging - what are we measuring?
  - e.g. pseudoscalar “charmonium” interpolating field  $\bar{c}\gamma^5 c$
  - at large Euclidean times (separations) should get the **lightest QCD eigenstate** that is pseudoscalar and isosinglet (and has some amplitude to be  $\bar{c}c$ )
  - this is the  $\eta$ -meson (via its hidden charm component) !
- separating properties of charmonium from the light-quark stuff is tricky
  - see recent work by **DeTar & Levkova**
  - non-trivial modeling of the correlator:

$$D(p^2) = \left[ \frac{f(p^2)}{p^2 + m_{\bar{c}c}^2} + \frac{g(p^2)}{p^2 + m_{\bar{c}c^*}^2} \right]^2 \frac{1}{p^2 + m_l^2}$$

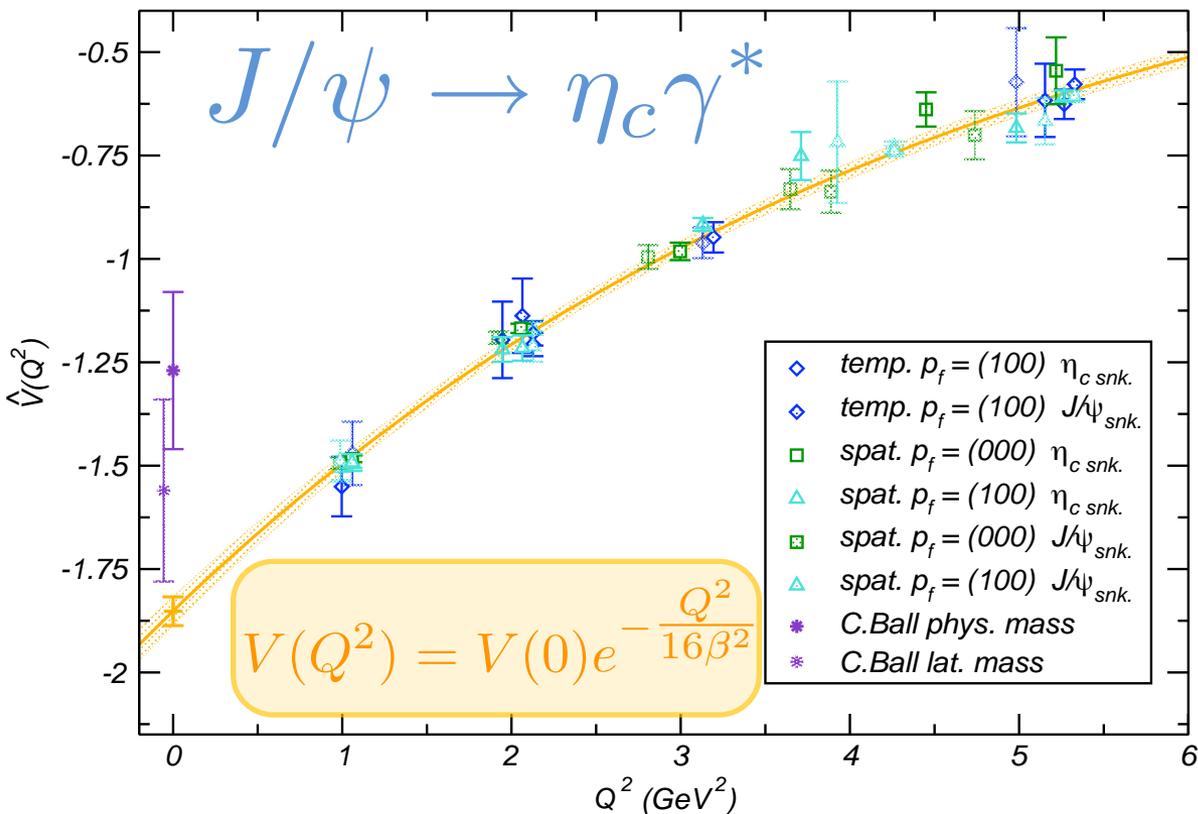


# radiative transitions

- first computation (quenched)
- no disconnected diagrams



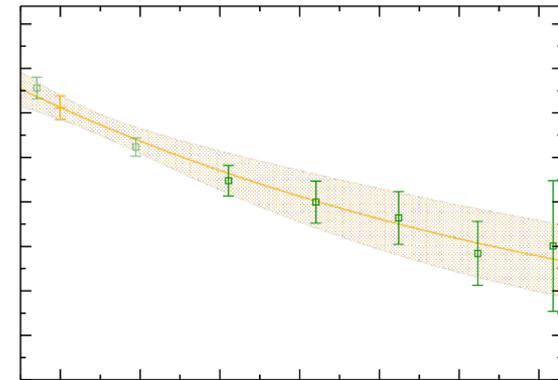
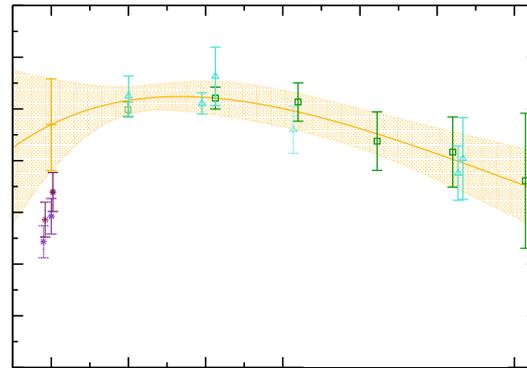
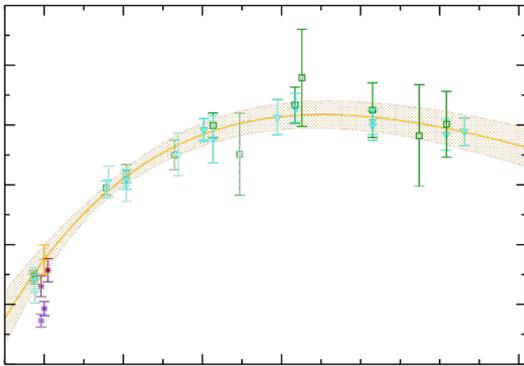
$J/\psi \rightarrow \eta_c \gamma^*$



# $1P \rightarrow 1S$ transitions

- fit form inspired by potential models with spin-dependent corrections

$$E_1(Q^2) = E_1(0) \left( 1 + \frac{Q^2}{\rho^2} \right) e^{-\frac{Q^2}{16\beta^2}}$$



$$\chi_{c0} \rightarrow J/\psi \gamma E1$$

$$\beta = 542(35) \text{ MeV}$$

$$\rho = 1.08(13) \text{ GeV}$$

$$\chi_{c1} \rightarrow J/\psi \gamma E1$$

$$\beta = 555(113) \text{ MeV}$$

$$\rho = 1.65(59) \text{ GeV}$$

$$h_c \rightarrow \eta_c \gamma E1$$

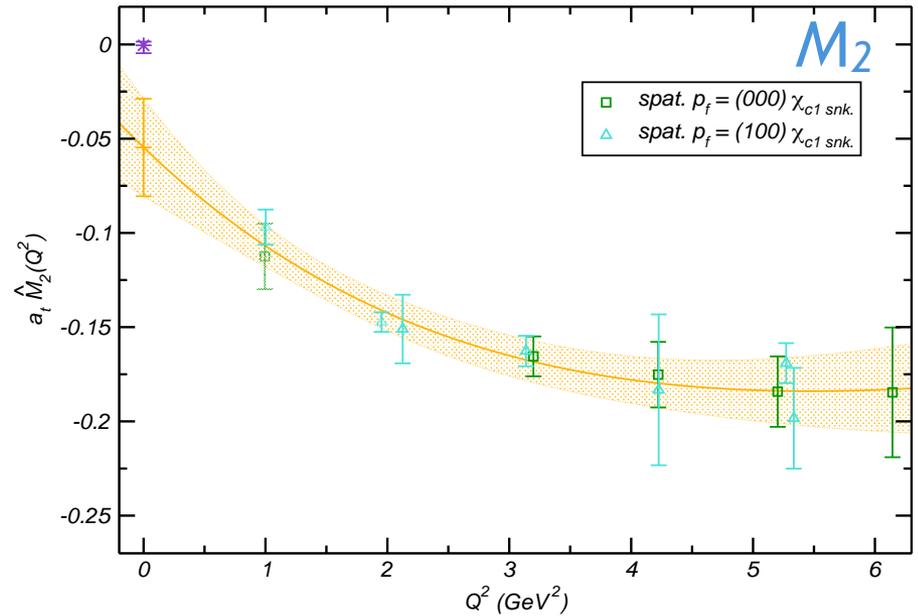
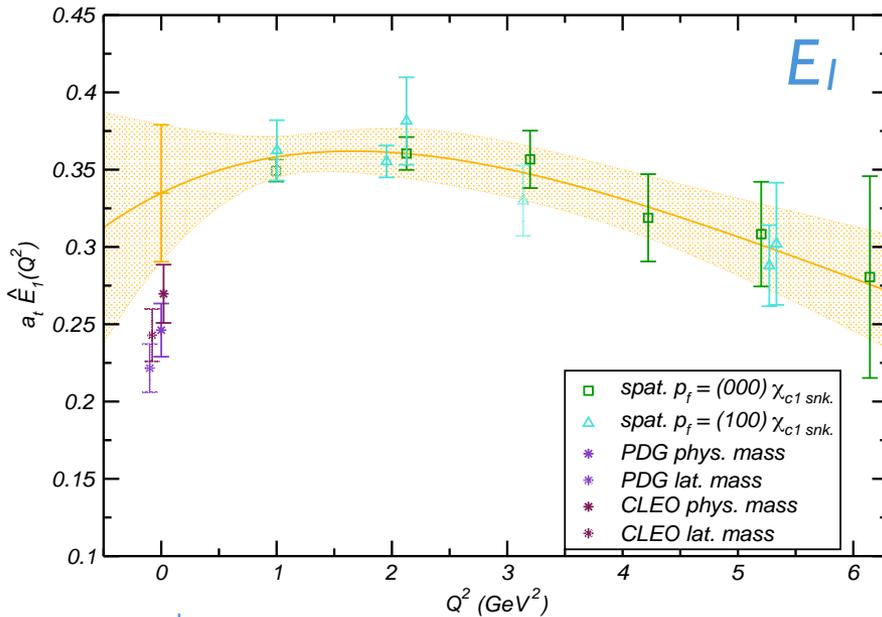
$$\beta = 689(133) \text{ MeV}$$

$$\rho \rightarrow \infty$$

simplest quark model has all  $\beta$  equal and  $\rho(\chi_{c0}) = 2\beta$ ,  $\rho(\chi_{c1}) = \sqrt{2} \cdot \rho(\chi_{c0})$ ,  $\rho(h_c) \rightarrow \infty$

# $\chi_{c1} \rightarrow J/\psi \gamma$ transition

- automatically access both the electric dipole and the magnetic quadrupole transitions



$$\left. \frac{M_2}{E_1} \right|_{\text{quen. lat}} = -0.2(1)$$

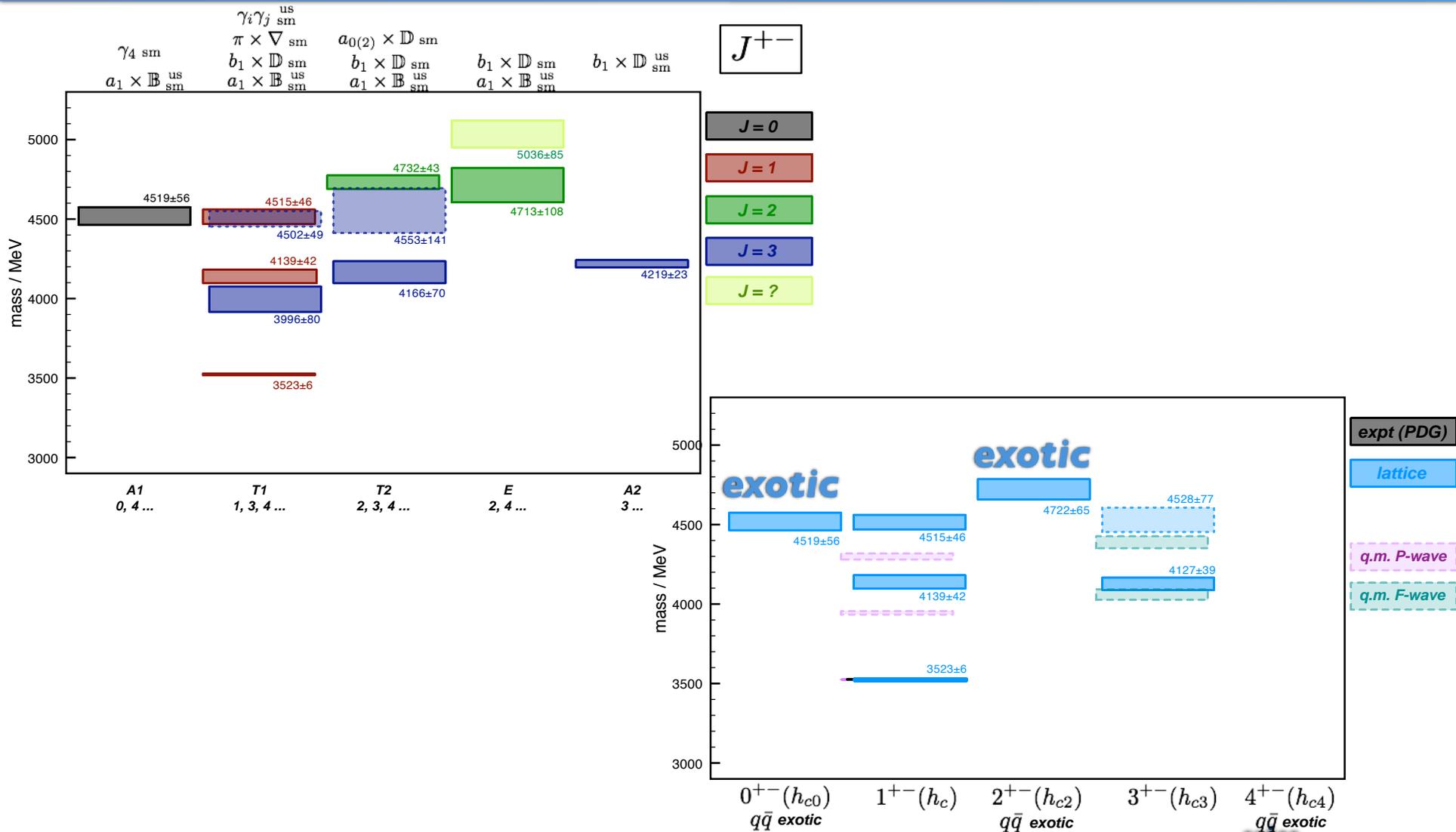
less theoretical motivation  
for the fit-form here

$$\left. \frac{M_2}{E_1} \right|_{\text{expt}} = -0.002 \left( \begin{smallmatrix} +8 \\ -17 \end{smallmatrix} \right)$$

# exotic quantum numbers ?

- $0^{--}, 0^{+-}, 1^{-+}, 2^{+-}$  cannot be constructed from a Fock state
- hence 'exotic' - no experimental charmonium candidates (to my knowledge)
- just build an operator with these quantum numbers!
- actually not quite as simple as it appears - lattice symmetry is not continuum rotations, discrete cubic rotations
  - $A_1 - 0, 4...$
  - $T_1 - 1, 3, 4...$
  - $T_2 - 2, 3, 4...$
  - $E - 2, 4...$
  - $A_2 - 3...$
- so the  $0^{+-}, 2^{+-}$  are probably straightforward
- but  $1^{-+}$  could be confused with a non-exotic  $4^{+-}$

# exotic quantum number?



# exotic quantum number?

$$\begin{array}{cccc}
 \gamma_5^{\text{us}} & \gamma_4 \gamma_5^{\text{us}} & a_0(2) \times \nabla_{\text{sm}} & \pi \times \mathbb{D}_{\text{sm}} \\
 b_1 \times \nabla_{\text{sm}}^{\text{us}} & b_1 \times \nabla_{\text{sm}} & b_1 \times \nabla_{\text{sm}} & b_1 \times \nabla_{\text{sm}} \\
 \rho \times \mathbb{B}_{\text{sm}}^{\text{us}} & \rho \times \mathbb{B}_{\text{sm}}^{\text{us}} & \rho \times \mathbb{B}_{\text{sm}} & \rho \times \mathbb{B}_{\text{sm}}^{\text{us}}
 \end{array}$$

$$J^{-+}$$

